

### Question

A viscous fluid has constant density  $\rho$  and constant kinematic viscosity  $\nu$ . Non-dimensionalise the steady Navier-Stokes equations with suitable scalings to show that, when the Reynolds number  $Re$  satisfies

$$Re \ll 1,$$

the non-dimensional equations of motion become, to leading order, the “slow-flow” equations

$$\begin{aligned}\nabla p &= \nabla^2 \underline{q} \\ \nabla \cdot \underline{q} &= 0\end{aligned}$$

where  $p$  and  $\underline{q}$  denote the fluid pressure and velocity respectively. Give an example of a fluid mechanics scenario where such a “slow-flow” approximation would apply.

The fluid is contained between two plates. Cartesian coordinates  $(x, z)$  are used. The plates are situated at  $z = 0$  and  $z = \delta h(x)$  in dimensionless variables where  $h(x)$  may be regarded as a given function and  $\delta \ll 1$ . Show how, by further scaling the slow flow equations according to

$$\begin{aligned}z &= \delta Z \\ w &= \delta W \\ p &= \frac{P}{\delta^2},\end{aligned}$$

the two-dimensional “lubrication theory” equations

$$\begin{aligned}P_x &= u_{ZZ} \\ P_Z &= 0 \\ u_x + W_Z &= 0\end{aligned}$$

may be derived. Given the conditions under which this limit is valid.

The top plate is held fixed, and the bottom plate is moved with a non-dimensional speed of 1 in the positive  $x$ -direction. Show that the pressure satisfies Reynolds’ equation

$$\frac{d}{dx} \left( \frac{h^3}{6} \frac{dP}{dx} \right) - \frac{dh}{dx} = 0$$

and give suitable boundary conditions for this differential equation. Suppose now that  $h(x) = a + bx$  where  $a$  and  $b$  are constants. WITHOUT SOLVING

the equation comment briefly on whether  $b$  should be positive or negative if the top plate is to support a load.

**Answer**

$$\text{We have } \left. \begin{aligned} (\underline{q} \cdot \nabla) \underline{q} &= -\frac{1}{\rho} \nabla p + \nu \nabla^2 \underline{q} \\ \text{div}(\underline{q}) &= 0 \end{aligned} \right\}$$

Since slow flow, we set  $\underline{x} = L\bar{x}$ ,  $p = (\mu U/L)\bar{p}$ ,  $\underline{q} = U\bar{q}$ ,  $\Rightarrow$

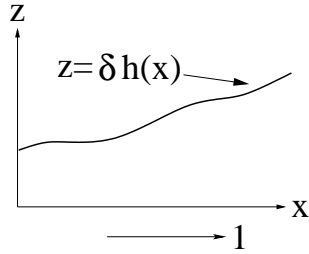
$$\left. \begin{aligned} \frac{U^2}{L} (\bar{q} \cdot \nabla) \bar{q} &= -\frac{1}{\rho L} \bar{\nabla} \bar{p} \left( \frac{\mu U}{L} \right) + \frac{\nu U}{L^2} \bar{\nabla}^2 \bar{q} \\ \bar{\nabla} \cdot \bar{q} &= 0 \end{aligned} \right\}$$

Multiplying the momentum equation by  $\nu U/L^2$

$$\Rightarrow \frac{LU}{\nu} (\bar{q} \cdot \bar{\nabla}) \bar{q} = -\bar{\nabla} \bar{p} + \bar{\nabla}^2 \bar{q}, \text{ but } \frac{LU}{\nu} = Re$$

Thus  $Re(\bar{q} \cdot \nabla) \bar{q} = -\bar{\nabla} \bar{p} + \bar{\nabla}^2 \bar{q}$ ,  $\Rightarrow$  for  $Re \ll 1$  we get, to lowest order,  $\bar{\nabla} \bar{p} = \bar{\nabla}^2 \bar{q}$ ,  $\bar{\nabla} \cdot \bar{q} = 0$

Any decent flow example will do, eg sperm swimming, treacle flowing, tar running down a telegraph pole.



We have (drop bars)

$$\left. \begin{aligned} p_x &= u_{xx} + u_{zz} \\ p_z &= w_{xx} + w_{zz} \\ u_x + w_z &= 0 \end{aligned} \right\}$$

So further scale  $\frac{1}{\delta^2} P$ ,  $w = \delta W$ ,  $z = \delta Z$ .

$$\Rightarrow \left. \begin{aligned} \frac{1}{\delta^2} P_x &= u_{xx} + \frac{1}{\delta^2} U_{ZZ} \\ \frac{1}{\delta} P_Z &= \delta W_{xx} + \frac{1}{\delta} W_{ZZ} \\ u_x + W_Z &= 0 \end{aligned} \right\}$$

Now to leading order as  $\delta \rightarrow 0$ , we must get

$$P_x = u_{ZZ}, \quad P_Z = 0, \quad u_x + W_Z = 0$$

Need  $\delta \ll 1$ ,  $\delta^2 Re \ll 1$ .

Now the top plate is fixed and the bottom is moved with speed 1.

$$\text{we have } u_{ZZ} = P_x \Rightarrow u = \frac{Z^2 P_x}{2} + AZ + B$$

Now  $u = 0$  on  $Z = h$ ,  $U = 1$  on  $Z = 0$

$$\Rightarrow B = 1, \quad o = \frac{h^2 P_x}{2} + Ah + 1$$

and so

$$\begin{aligned}
 u &= 1 + \frac{Z^2 P_x}{2} + Z \left( -\frac{1}{h} - \frac{h P_x}{2} \right) \\
 &= \frac{Z(Z-h)P_x}{2} + \left( 1 - \frac{Z}{h} \right) \\
 \int_h^0 (u_x + W_Z) dZ &= 0 \Rightarrow \int_0^h u_x dZ = 0 \quad (\text{since } W = 0 \text{ on } z = 0, h) \\
 \Rightarrow \frac{\partial}{\partial x} \int_0^h u dZ &= 0 \quad (\text{since } u = 0 \text{ on } Z = h) \\
 \frac{d}{dx} \left[ \frac{Z^3}{6} P_x - \frac{Z^2 h}{4} P_x + Z - \frac{Z^2}{2h} \right]_0^h &\Rightarrow \frac{d}{dx} \left[ -\frac{h^3}{12} P_x + \frac{h}{2} \right] = 0 \\
 \Rightarrow \frac{d}{dx} \left( \frac{h^3}{6} P_x \right) - \frac{dh}{dx} &= 0
 \end{aligned}$$

We impose  $P = 0$  at  $x = 0, 1$

When  $h = a + bx$  need  $b < 0$  for  $P > 0$  so that fluid is being forced into a NARROWING gap.